

Polarization splitting and rotating through adiabatic transitions

M. R. Watts, G. Gorni, M. Cherchi, H. A. Haus

*Research Laboratory of Electronics, Department of Electrical Engineering and
Computer Science, Massachusetts Institute of Technology, MIT 36-355, 50
Vassar St., Cambridge, MA 02139*

Tel (617)-253-5481, Fax (617)-253-9611, mwatts@mit.edu

Abstract: A polarization splitter and rotator that is compatible with optical integration is presented. In contrast to prior efforts to achieve this functionality, the adiabatic nature of this approach allows for nearly ideal performance over the 1.5-1.6 μm regime.

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OCIS codes: (130.3120) Integrated optics devices; (130.2790) Guided Waves

High index contrast waveguides allow for small bend radii and thus enable dense integration of optical devices [3]. In high index-contrast systems it is difficult, if not impossible, to maintain degeneracy of the quasi-TE and quasi-TM modes across each and every structure on the chip thus causing polarization dependent performance. One approach to eliminate polarization sensitivity is to split the random polarization state emanating from the fiber into two orthogonal polarization states, rotate one of the states, and then perform identical operations on each in parallel [6]. In this way, the design and fabrication need only be optimized for a single polarization.

For dense integration and elimination of the group delay between the paths, it is desirable to achieve both the splitting and rotation of polarization states in an integrated manner. In this configuration, the polarization splitting and rotating is done at the chip input, and thus must perform effectively across all optical channels propagating on the fiber. Any such approach must therefore possess low loss across a very broad optical spectrum.

We consider rotation of the polarization. Many approaches to polarization rotation have been proposed [2, 4, 5, 6]. In each, the orthogonally polarized quasi-TE and quasi-TM modes of the guide are in one way or another coupled in order to invoke rotation. Coupling between orthogonally polarized modes is inherently difficult since the coupling must be achieved through minor field components. Thus, such structures are often complicated and difficult to fabricate. Moreover, coupled modes must be phase-matched, have precisely tuned coupling coefficients, and suffer from an inherent wavelength sensitivity. We propose a structure that seeks to transition the modes without invoking coupling. The approach is based on an adiabatic transition from a vertically aligned rectangular waveguide to a horizontally aligned rectangular waveguide. Although a pure rotation would be ideal, standard fabrication techniques dictate a layered structure. We consider a three layered approach with the core divided equally amongst the three-layers (Fig. 1a). The transition is made in a linear fashion with cross-sectional area removed from the top and bottom layers and added to the middle layer. By starting with a large aspect ratio rectangular waveguide we ensure the quasi-TE and quasi-TM modes propagate with greatly differing velocities and therefore dephase before they have a chance to exchange substantial power. Thus, despite the coupling introduced by the transition, the phase mismatch between the quasi-TE and quasi-TM modes mitigates coupling between them enabling rotation without mode mixing.

The three-layered approach adopted for rotating the polarization allows for a natural approach to adiabatically split the polarization states within the same three layers. To split the polarizations, we begin with a horizontally oriented rectangular waveguide centrally intersected by an identical vertically oriented guide therein forming a cross structure (Fig. 1b). The quasi-TE and quasi-TM modes of the structure are degenerate. Despite this degeneracy upon slowly separating the guides there is no mechanism for the two modes to couple since the modes remain orthogonal throughout the length of the structure. These primary modes, however, may couple to secondary modes which appear as the two structures become distinct and lack the appropriate field symmetry to prevent coupling with their like polarized counterparts. However, the mode sets are not phase-matched and thus little power transfer occurs between the modes when the transition is made to be sufficiently adiabatic.

Verification of the performance of these devices is achieved through the use of the three-dimensional mode

scattering formulation of [1]. The structures are those of Fig. 1a and Fig. 1b with $a = 0.25\mu\text{m}$, $b = 0.75\mu\text{m}$, a core index of 2.2 a cladding index of 1.445 with device lengths of $400\mu\text{m}$ and $150\mu\text{m}$ for the rotator and splitter, respectively. Scans of the transmission as a function of wavelength are presented in Fig. 2a and Fig. 2b. In each case greater than 99% of the power is transferred successfully across the entire 1.5-1.6 μm band with greater than 25dB of suppression of the unwanted polarization states.

Variations of these basic device designs will be presented along with an integrated mechanism for coupling to a standard SMF fiber.

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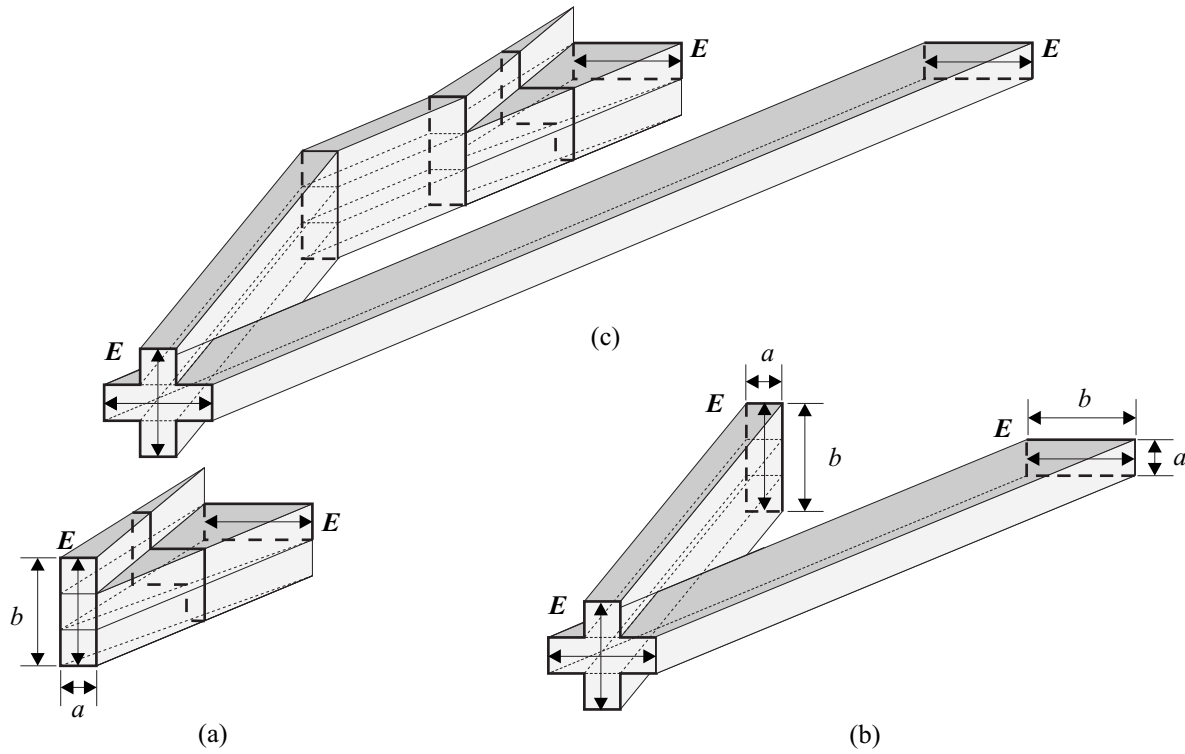


Figure 1. (a) Polarization rotator, (b) polarization splitter, and (c) integrated polarization splitter and rotator.

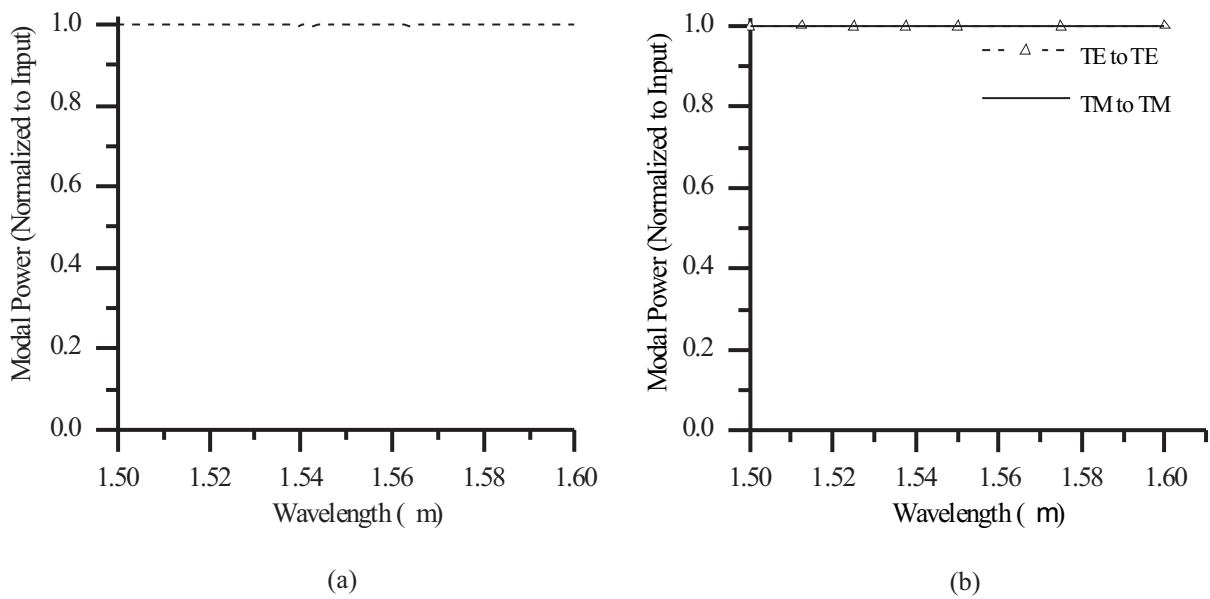


Figure 2. Transmission for (a) a 400 μm long polarization rotator and (b) a 150 μm long polarization splitter with a 1 μm separation of the core walls at the output. In each case $a = 0.25 \mu\text{m}$, $b = 0.75 \mu\text{m}$, and the core and cladding indices are 2.2 and 1.445.